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# Solution of Heat and Wave Equations using Mahgoub Adomian Decomposition Method

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**ABSTRACT**: In this paper, Mahgoub Adomian decomposition method (MADM), to handle the wave and heat equations, is introduced. The efficiency of the present method will be shown by applying the procedure on four examples.

*Keywords* – *Mahgoub Adomian decomposition method, Mahjoub transform, Partial Differential Equations, Heat equation, Wave equation.* 

# INTRODUCTION

An *n*-th order PDE for  $u(x_1, x_2, x_3, ..., x_n)$ ,  $n \ge 2$  is a relation of the form  $F(\vec{x}, u, Du, D^2u, ..., D^nu) = 0$ ,

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where F is neither linear or nonlinear function.

The focus will be on two main types of PDEs as well, namely wave and heat equations. Now, we are going to define these equations as

$$\frac{\partial^2 u}{\partial t^2} = c^2 \left( \frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_n^2} \right)$$
(2)

$$\frac{\partial u}{\partial t^2} = c^2 \left( \frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_n^2} \right)$$
(3)

respectively. The wave and heat equations have many possible applications in mathematics, physics and engineering (see [1-6]).

In 2016, Mahgoub [7] define the function A for  $t \ge 0$ 

$$A = \left\{ u(t) : \exists M, k_1, k_2 > 0, |u(t)| < M e^{\frac{|t|}{k_j}} \right\},\tag{4}$$

where M,  $k_1$ ,  $k_2$  are fixed and M is a finite.

The operator  $M{u(t)}$  may be expanded as

$$M\{u(t)\} = H(v) = v \int_0^\infty u(t) e^{-vt} dt, \quad k_1 \le v \le k_2.$$
(5)

Many linear and nonlinear problems like ODEs, PDEs, Integral equations and integro- differential equations were solved using Mahgoub transform and ADM(see [8 - 13]).

In present paper, we have application MADM for solving wave and heat equations.

# II. APPLICATION OF THE METHOD

The general inhomogeneous form of PDE is

$$Lu(x,t) + Ru(x,t) + Nu(x,t) = h(x,t)$$
(6)

with

u(x, 0) = f(x)  $u_t(x, 0) = g(x)$ 

where  $L = \frac{\partial^2}{\partial t^{2'}}$  with inverse  $L_u^{-1} = \int_0^t \int_0^t (.) dt dt$ , *R* is linear operator and *N* is nonlinear operator. Now, we rewrite (6) as

$$Lu(x,t) = h(x,t) - Ru(x,t) - Nu(x,t)$$
Take Mahgoub transform to both sides, we have
(7)

$$M\{Lu(x,t)\} = M\{h(x,t) - Ru(x,t) - Nu(x,t)\}$$
(8)

Linearity of 
$$M\{.\}$$
 yields

$$M\{Lu(x,t)\} = M\{h(x,t)\} - M\{Ru(x,t)\} - M\{Nu(x,t)\}$$
Solving for (9), we get
(9)

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 $v^{2}M\{u(x,t)\} - u(x,0)v^{2} - vu_{t}(x,0) = M\{h(x,t)\} - M\{Ru(x,t)\} - M\{Nu(x,t)\}$ (10)By substituting u(x, 0) and  $u_t(x, 0)$  into (10), we obtain

$$v^2 M\{u(x,t)\} - f(x)v^2 - vg(x) = M\{h(x,t)\} - M\{Ru(x,t)\} - M\{Nu(x,t)\}$$
(11)  
This leads to

$$M\{u(x,t)\} = f(x) + \frac{1}{v}g(x) + \frac{1}{v^2}M\{h(x,t)\} - \frac{1}{v^2}M\{Ru(x,t)\} - \frac{1}{v^2}M\{Nu(x,t)\}$$
(12)

Replacing  $u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$  and  $Nu(x,t) = \sum_{n=0}^{\infty} A_n$  in (12), we have  $M\{\sum_{n=0}^{\infty} u_n(x,t)\} = f(x) + \frac{1}{v}g(x) + \frac{1}{v^2}M\{h(x,t)\} - \frac{1}{v^2}M\{R\sum_{n=0}^{\infty} u_n(x,t)\} - \frac{1}{v^2}M\{\sum_{n=0}^{\infty} A_n\}$ (13) where

$$A_{n} = \frac{1}{n!} \frac{d^{n}}{d\lambda^{n}} [\sum_{i=0}^{\infty} \lambda^{n} N u_{i}]_{\lambda=0}, \quad m = 0, 1, \dots$$
(14)

Therefore, we have

$$M\{u_0(x,t)\} = f(x) + \frac{1}{v}g(x) + \frac{1}{v^2}M\{h(x,t)\}$$
  

$$M\{u_{n+1}(x,t)\} = -\frac{1}{v^2}(M\{R\sum_{n=0}^{\infty}u_n(x,t) + M\{\sum_{n=0}^{\infty}A_n\}\}), n \ge 0$$
(15)  

$$M^{-1} \text{ to both sides yields}$$

Applying  $M^{-1}$  to both sides, yields

$$u_0(x,t) = M^{-1}\{f(x)\} + tg(x) + M^{-1}\left\{\frac{1}{v^2} M\{h(x,t)\}\right\}$$
$$u_{n+1}(x,t) = -M^{-1}\left\{\frac{1}{v^2}(M\{R\sum_{n=0}^{\infty} u_n(x,t) + M\{\sum_{n=0}^{\infty} A_n\}\})\right\}, \ n \ge 0$$
(16)

#### III. **APPLICATIONS AND RESULTS**

**Example 1:** Consider the wave equation [5]

$$u_{tt} = \frac{1}{2} x^2 u_{xx}, \qquad (17)$$

Subject to u(x, 0) = x,  $u_x(x, 0) = x^2$ . We use  $M\{.\}$  to both sides, yields

$$M\{u_{tt}\} = M\left\{\frac{1}{2}x^2u_{xx}\right\},\tag{18}$$

which gives

$$v^{2}M\{u(x,t)\} - u(x,0)v - vu_{t}(x,0) = M\left\{\frac{1}{2}x^{2}u_{xx}\right\}$$
By substituting  $u(x,0) = x$ ,  $u_{x}(x,0) = x^{2}$  into (19), we obtain
$$(19)$$

$$v^{2}M\{u(x,t)\} - xv^{2} - vx^{2} = M\left\{\frac{1}{2}x^{2}u_{xx}\right\}$$
(20)  
Replacing  $u(x,t) = \sum_{n=0}^{\infty} u_{n}(x,t)$  in (20), we have

$$M\{\sum_{n=0}^{\infty} u_n(x,t)\} = x + \frac{1}{v}x^2 + \frac{1}{v^2}M\left\{\frac{1}{2}x^2u_{nxx}\right\}, \ n \ge 0$$
(21)

This leads to

Is leads to  

$$M\{u_0(x,t)\} = x + \frac{1}{v}x^2$$

$$M\{u_{n+1}(x,t)\} = \frac{1}{v^2}M\{\frac{1}{2}x^2u_{n_{xx}}\}, n \ge 0$$
(22)
plying  $M^{-1}$  to both sides yields

Applying  $M^{-1}$  to both sides, yields

$$u_{0}(x,t) = x + tx^{2}$$

$$u_{n+1}(x,t) = M^{-1} \left\{ \frac{1}{v^{2}} M \left\{ \frac{1}{2} x^{2} u_{n_{XX}} \right\} \right\}, \ n \ge 0$$
(23)

The solutions of (23) are given by  $u_0(x,t) = x + tx^2$ 

$$\begin{aligned} u_1(x,t) &= M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 u_{0_{xx}} \right\} \right\} = M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 (2t) \right\} \right\} = x^2 M^{-1} \left\{ \frac{1}{v^3} \right\} = \frac{1}{3!} t^3 x^2 \\ u_2(x,t) &= M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 u_{1_{xx}} \right\} \right\} = M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 \left( \frac{2}{3!} t^3 \right) \right\} \right\} = \left( \frac{x^2}{3!} \right) M^{-1} \left\{ \frac{1}{v^5} \right\} = \frac{x^2 t^5}{5!} \\ u_3(x,t) &= M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 u_{2_{xx}} \right\} \right\} = M^{-1} \left\{ \frac{1}{v^2} M \left\{ \frac{1}{2} x^2 \left( \frac{2}{5!} t^5 \right) \right\} \right\} = \frac{1}{7!} x^2 M^{-1} \left\{ \left\{ \frac{1}{v^7} \right\} \right\} = \frac{x^2 t^7}{7!} \end{aligned}$$

Thus

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t) = x + tx^2 + \frac{x^2 t^3}{3!} + \frac{x^2 t^5}{5!} + \cdots$$
  
=  $x + x^2 \left( t + \frac{t^3}{3!} + \frac{t^5}{5!} + \frac{t^7}{7!} + \cdots \right)$   
=  $x + x^2 \sinh t$  (24)

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Example 2: Consider the Heat equation [4]	
$u_t = u_{xx}$ ,	(25)
Subject to $u(x, 0) = ae^{kx}$ , where a and k are constants.	
We use $M\{.\}$ to both sides, yields	
$M\{u_t\} = M\{u_{xx}\},$	(26)
which gives	

$$vM\{u(x,t)\} - vu(x,0) = M\{u_{xx}\}$$
By substituting  $u(x,0) = ae^{kx}$  into (27), we obtain
$$(27)$$

$$vM\{u(x,t)\} - ae^{kx}v = M\{u_{xx}\} \Rightarrow M\{y(x,t)\} = ae^{kx} + \frac{1}{v}M\{u_{xx}\}$$
(28)

Replacing  $u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$  in (28), we have  $M\{\sum_{n=0}^{\infty} u_n(x,t)\} = ae^{kx} + \frac{1}{v}M\{u_{n_{xx}}\}$ This leads to (29)

$$M\{u_0(x,t)\} = ae^{kx}$$

$$M\{u_{n+1}(x,t)\} = \frac{1}{v}M\{u_{n_{xx}}\}, n \ge 0$$
(30)

Applying  $M^{-1}$  to both sides, yields  $u_0(x,t) = ae^{kx}$ 

$$u_{n+1}(x,t) = M^{-1} \left\{ \frac{1}{v} M\{u_{n_{xx}}\} \right\}, \ n \ge 0$$
(31)

The solutions of (31) are given by  $u_0(x,t) = ae^{kx}$ 

$$\begin{aligned} u_1(x,t) &= M^{-1} \left\{ \frac{1}{v} M\{u_{0xx}\} \right\} = M^{-1} \left\{ \frac{1}{v} M\{ak^2 e^{kx}\} \right\} = ak^2 e^{kx} M^{-1} \left\{ \frac{1}{v} \right\} = ak^2 t e^{kx} \\ u_2(x,t) &= M^{-1} \left\{ \frac{1}{v} M\{u_{1xx}\} \right\} = M^{-1} \left\{ \frac{1}{v} M\{ak^4 t e^{kx}\} \right\} = M^{-1} \left\{ ak^4 e^{kx} \frac{1}{v} M\{t\} \right\} \\ &= ak^4 e^{kx} M^{-1} \left\{ \frac{1}{v^2} \right\} = ak^4 \frac{t^2}{2!} e^{kx} \\ u_3(x,t) &= M^{-1} \left\{ \frac{1}{v} M\{u_{2xx}\} \right\} = M^{-1} \left\{ \frac{1}{v} M\{ak^6 \frac{t^2}{2!} e^{kx} \right\} = M^{-1} \left\{ \frac{ak^6 e^{kx}}{2!} \frac{1}{v} M\{t^2\} \right\} \\ &= \frac{ak^6 e^{kx}}{2!} M^{-1} \left\{ \frac{2}{v^3} \right\} = ak^6 \frac{t^3}{3!} e^{kx} \end{aligned}$$

Thus

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t) = ae^{kx} + ak^2 te^{kx} + ak^4 \frac{t^2}{2!} e^{kx} + \cdots$$
  
=  $ae^{kx} \left( 1 + k^2 t + k^4 \frac{t^2}{2!} + k^6 \frac{t^3}{3!} + \cdots \right)$   
=  $ae^{kx} e^{k^2 t}$   
=  $ae^{kx+k^2 t}$  (32)

**Example 3:** The following nonlinear heat equation [4]

 $u_t + uu_x = u_{xx}$  , (33)with u(x, 0) = 2x.

Mahgoub transform of (33) is

$$M\{u_t\} = M\{u_{xx}\} - M\{uu_x\},$$
(34)

Therefore

$$vM\{u(x,t)\} - vu(x,0) = M\{u_{xx}\} - M\{uu_x\}$$
(35)

By substituting u(x, 0) = 2x into (35), we find that

$$M\{u(x,t)\} = 2x + \frac{1}{v}M\{u_{xx}\} - \frac{1}{v}M\{uu_x\}$$
(36)

Replacing  $u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$  and  $Nu(x,t) = \sum_{n=0}^{\infty} A_n$  in (36) and then applying  $M^{-1}$ , yields  $u_0(x,t) = 2x$ 

$$u_{n+1}(x,t) = M^{-1} \left\{ \frac{1}{v} \left( M\{u_{n_{xx}}\} - M\{A_n\} \right) \right\}, \ n \ge 0$$
(37)
the solutions of (37) are

Therefore the (37) $u_0(x,$ 

$$\begin{aligned} &(t) = 2x \\ &u_1(x,t) = M^{-1} \left\{ \frac{1}{v} \left( M\{u_{0_{xx}}\} - M\{A_0\} \right) \right\} = M^{-1} \left\{ \frac{1}{v} \left( M\{0\} - M\{u_0u_x\} \right) \right\} = M^{-1} \left\{ -\frac{1}{v} 4x \right\} \end{aligned}$$

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$$= -4xM\left\{\frac{1}{v}\right\} = -4xt$$

$$u_{2}(x,t) = M^{-1}\left(\frac{1}{v}\left(M\{u_{1_{xx}}\} - M\{A_{1}\}\right)\right) = M^{-1}\left\{\frac{1}{v}\left(M\{0\} - M\{u_{0_{x}}u_{1} + u_{1_{x}}u_{0}\}\right)\right\}$$

$$= M^{-1}\left\{-\frac{1}{v}M\{-16xt\}\right\} = 16xM^{-1}\left\{\frac{1}{v}M\{t\}\right\} = 16xM^{-1}\left\{\frac{1}{v^{2}}\right\} = 16x\frac{t^{2}}{2!}$$

$$u_{3}(x,t) = M^{-1}\left\{\frac{1}{v}\left(M\{u_{2_{xx}}\} - M\{A_{2}\}\right)\right\} = -96x\frac{t^{3}}{3!},$$
Then, the solution is

and so on. Then, the solution is

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t) = -4xt + 16x \frac{t^2}{2!} - 96x \frac{t^3}{3!} + \cdots$$
  
=  $\frac{2x}{1+2t}$ . (38)

**Example 4:** Solve the nonlinear wave equation [5]

$$u_t + u_x^2 = 0,$$
Subject to  $u(x, 0) = -x^2.$ 
Take  $M\{.\}$  to both sides, we have
$$(39)$$

$$M\{u_t\} = -M\{u_x^2\},$$
(40)

And

$$vM\{u(x,t)\} - vu(x,0) = -M\{u_x^2\}$$
(41)

Thus

$$M\{u(x,t)\} = -x^2 - \frac{1}{v} M\{u_x^2\}$$
(42)

Using  $u(x,t) = \sum_{n=0}^{\infty} u_n(x,t)$  and  $Nu(x,t) = \sum_{n=0}^{\infty} A_n$  in (42) and then taking  $M^{-1}$ , yields  $u_0(x,t) = -x^2$ 

$$u_{n+1}(x,t) = -M^{-1} \left\{ \frac{1}{v} \left( M\{A_n\} \right) \right\}, \ n \ge 0$$
(43)

,

. .

Therefore the solutions of (43) are given

 $u_0(x,t) = -x^2$ 

$$\begin{aligned} u_1(x,t) &= -M^{-1} \left\{ \frac{1}{v} M(A_0) \right\} = -M^{-1} \left\{ \frac{1}{v} M\{u_{0_x}^2\} \right\} = -M^{-1} \left\{ \frac{4x^2}{v} \right\} \\ &= -4x^2 M \left\{ \frac{1}{v} \right\} = -4x^2 t \\ u_2(x,t) &= -M^{-1} \left\{ \frac{1}{v} M\{A_1\} \right\} = -M^{-1} \left\{ \frac{1}{v} M\{2u_{0_x}u_{1_x}\} \right\} = -2M^{-1} \left\{ \frac{1}{v} M\{16x^2t\} \right\} \\ &= -32x^2 M^{-1} \left\{ \frac{1}{v} M\{t\} \right\} = -32x^2 M^{-1} \left\{ \frac{1}{v^2} \right\} = -32x^2 \frac{t^2}{2!} = -16x^2 t^2 \\ u_3(x,t) &= -M^{-1} \left\{ \frac{1}{v} M\{A_2\} \right\} = -64x^2 t^3 , \\ u_m(x,t) &= -M^{-1} \left\{ \frac{1}{v} M\{A_{m-1}\} \right\} = -4^m x^2 t^m, \ m = ,1,2, \dots \end{aligned}$$
 lution is

Then, the solution i

$$u(x,t) = \sum_{n=0}^{\infty} u_n(x,t) = -x^2 - 4x^2t - 16x^2t^2 + \cdots$$
  
=  $\frac{x^2}{1+4t}$ . (44)

## IV. CONCLUSION

Mahgoub Adomian decomposition method to the wave and heat equations was introduced. The efficiency of the method has been shown by applying the procedure on some examples.

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